

# Search for Galactic dark matter substructures with Cherenkov telescopes

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Weakly interacting massive dark matter (DM) particles are expected to self-annihilate or decay, generating high-energy photons in these processes. This establishes the possibility for indirect detection of DM by  $\gamma$ -ray telescopes. For probing the secondary products of DM, accurate knowledge about the DM density distribution in potential astrophysical targets is crucial. In this contribution, the prospects for the detection of subhalos in the Galactic DM halo with present and future imaging atmospheric Cherenkov telescopes (IACT) are investigated. The source count distribution and angular power spectra for  $\gamma$ -rays originating from annihilating DM in subhalos are calculated from N-body simulation results. To study the systematic uncertainties coming from the modeling of the DM density distribution, parameters describing the  $\gamma$ -ray yield from subhalos are varied in 16 benchmark models. We conclude that Galactic subhalos of annihilating DM are probably too faint to be a promising target for IACT observations, even with the prospective Cherenkov Telescope Array (CTA).

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## 1. Introduction

Revealing the nature of dark matter (DM) constitutes one of the most challenging tasks for astrophysics and cosmology. The most promising candidate to explain the numerous observational evidence for DM is a non standard model weakly interacting massive particle (WIMP), which self-annihilates into standard model particles [1]. As the WIMP mass is expected to be in the range of GeV to TeV, it is supposed to produce secondary  $\gamma$ -rays in the same energy regime. This opens the window for an indirect detection of DM by  $\gamma$ -ray telescopes. Various techniques exist for the detection of astrophysical  $\gamma$ -rays. This contribution is focused on imaging atmospheric Cherenkov telescopes (IACT), in particular the prospective Cherenkov Telescope Array (CTA) [2].

So far, no clear evidence for particle DM, neither direct nor indirect, has been found. Several targets have been considered for the indirect detection of DM, among them the Galactic center, neighboring dwarf spheroidal galaxies (dSphG), and galaxy clusters [3]. Extensive numerical simulations in the framework of  $\Lambda$ CDM cosmology have shown that DM is also expected to form density structures on sub-galactic scales [4, 5]. Such subhalos in our Galaxy may constitute an interesting additional target for indirect DM searches. First, they are expected to lack of any contaminating non-DM  $\gamma$ -ray emission, resulting into a signal-to-noise ratio comparable or better than for the dSphG observations. Second, as the flux of secondary  $\gamma$ -rays is expected to be proportional to the DM density squared, DM subhalos may cause a comparably high small-scale fluctuation of the isotropic  $\gamma$ -ray background (IGRB), detectable and distinguishable from other sources [6].

Many studies have been performed to assess the  $\gamma$ -ray fluxes from Galactic DM subhalos, by using the results of numerical simulations, e.g., [7, 8]. However, uncertainties about the spatial component of the signal, the *astrophysical J-factor*, still remain. The resolution of numerical N-body simulations is limited to larger subhalo masses than relevant for a detectable signal. Therefore, extensive extrapolations and model assumptions have to be applied over many orders of magnitude. In this contribution, these model systematic uncertainties are studied with the new version of the semi-analytic code CLUMPY [9, 10]. Four different N-body simulation fit results and extrapolations are varied, and an uncertainty band is derived for the *J*-factors and the fluctuations in the IGRB from DM subhalos in the Milky Way.

## 2. $\gamma$ -rays from Galactic dark matter subhalos

The  $\gamma$ -ray flux  $S$  from annihilating dark matter at redshift  $z = 0$  can be expressed as

$$S = \frac{1}{8\pi} \frac{\langle \sigma v \rangle}{m_\chi^2} \int_{E_\gamma} \sum_i b_i \left( \frac{dN_\gamma}{dE_\gamma} \right)_i \cdot \int_{\Delta\Omega} \int_{l.o.s.} \rho^2[r(l, \Omega)] dl d\Omega, \quad (2.1)$$

where  $m_\chi$  is the dark matter particle's mass,  $\langle \sigma v \rangle$  the velocity-averaged annihilation cross section,  $b_i$  is the branching ratio for annihilation channel  $i$ , and  $(dN_\gamma/dE_\gamma)_i$  the final-state  $\gamma$ -ray spectrum in this channel. The second term, the *J*-factor,  $J = \int \int \rho^2 dl d\Omega$ , is the integral of the squared DM density,  $\rho$ , over the line of sight,  $l$ , and  $\Delta\Omega = 2\pi(1 - \cos \theta)$ , the solid angle over the instrument's aperture radius,  $\theta$ . The flux  $S$  is proportional to  $J$ , which is independent of any particle physics model. Therefore, in the following all quantities are expressed in terms of  $J$  instead of  $S$ .

The mean density  $\bar{\rho}$  of the total Galactic DM host halo is modeled by an Einasto profile,

$$\bar{\rho}(r) = \bar{\rho}_s \exp \left( -\frac{2}{\alpha_E} \left[ \left( \frac{r}{r_s} \right)^{\alpha_E} - 1 \right] \right), \quad (2.2)$$

on the basis of a best fit to the Aquarius simulations presented in [8], with the scale radius  $r_s = 15.14 \text{ kpc}$  and  $\alpha_E = 0.17$ . The parameter  $1/\alpha_E$  is a measure of the steepness of the density profile close to  $r = 0$ , with  $r$  the distance from the halo center. Motivated by recent astronomical measurements [11], the halo is normalized by a local dark matter density of  $0.4 \text{ GeV cm}^{-3}$  at  $r_\odot = 8.0 \text{ kpc}$ . For a virial radius of  $260 \text{ kpc}$  of the host halo, this results in a total halo mass of  $1.1 \cdot 10^{12} M_\odot$ . This mass is in agreement with observations of stellar tracers [11] and is used in the following.

The simulations indicate that the distribution of subhalos in the host halo follows an anti-biased distribution with respect to the unclustered, smooth DM distribution [4, 5]. This implies that most of the clustered mass is located in the outer parts of the host halo. For the model variation, two parametrizations of the substructure distribution  $\bar{\rho}_{sub}$  within the host halo are compared: The Aquarius-simulation [4] suggests an Einasto profile for the subclustered mass with  $\alpha_E = 0.678$  and  $r_s = 199 \text{ kpc}$  (model EINASTO). The Via-Lactea II subhalos [5] are described by a slightly less anti-biased distribution, modeled by the number distribution  $N_{sub}(< r)$  given in [12] (model MADAU). All subhalos are described by an Einasto profile with the same  $\alpha_E$  as for the total halo.

The presence of subhalos with masses in the range  $[10^{-6} M_\odot, 10^{10} M_\odot]$  is assumed [13]<sup>1</sup>. The simulations [4, 5] indicate a power-law function  $dN_{sub}/dM_{sub} \sim M_{sub}^{-\alpha_m}$  of the subhalo distribution, with an index  $1.9 \lesssim \alpha_m \lesssim 2.0$  for the resolved subhalos with masses  $M_{sub} \gtrsim 10^5 M_\odot$ . The extrapolation of  $dN_{sub}/dM_{sub}$  down to the lowest subhalo masses is very sensitive to  $\alpha_m$ . Therefore, the results for the two indices  $\alpha_m = 1.9$  and  $\alpha_m = 2.0$  are compared. A fixed number of 150 massive dSphG-like subhalos in the mass range  $[10^8 M_\odot, 10^{10} M_\odot]$  is supposed, in agreement with [4].

An important property of a subhalo is the concentration  $c(M_{sub}) = R_\Delta/r_{s,sub}$ , which connects its scale radius  $r_{s,sub}$  with its virial size  $R_\Delta$  as a function of its mass  $M_{sub}$ . As the annihilation  $\gamma$ -ray emission is proportional to the density squared, a higher concentration significantly boosts the brightness of a subhalo. The rather optimistic (highly concentrated) model from [14] (model BULLOCK) is compared to the conservative one (flatter subhalo concentrations) from [13] (model SANCHEZ). Finally, the authors of [14] find a significant log-norm scattering  $\sigma_c = \Delta(\log c)$  around the mean concentrations  $\bar{c}(M_{sub})$ , which also results into a net boost of the signal. Their result for hosted subhalos,  $\sigma_c = 0.24$ , is compared to a more conservative assumption of  $\sigma_c = 0.14$ .

From the two values for each of the four varied quantities, 16 benchmark models are created, listed in Tab. 1. Model 9 is chosen to be the reference model with an average flux yield.

Model	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16
$\bar{\rho}_{sub}$	M	E	M	E	M	E	M	E	M	E	M	E	M	E	M	E
$\alpha_m$	1.9	1.9	2.0	2.0	1.9	1.9	2.0	2.0	1.9	1.9	2.0	2.0	1.9	1.9	2.0	2.0
$c(M_{sub})$	S	S	S	S	B	B	B	B	S	S	S	S	B	B	B	B
$\sigma_c$	0.14	0.14	0.14	0.14	0.14	0.14	0.14	0.14	0.24	0.24	0.24	0.24	0.24	0.24	0.24	0.24
$\alpha_J$	2.05	2.06	2.15	2.15	2.03	2.02	2.12	2.13	2.04	2.05	2.14	2.15	2.03	2.04	2.09	2.12

**Table 1:** Parameters for the 16 models investigated in this study. The reference model is marked in gray. S/B stands for the SANCHEZ/BULLOCK  $c(M_{sub})$  parametrization, E/M:  $\bar{\rho}_{sub}$  parametrization by an EINASTO/MADAU profile. Bottom row: Fit values for the exponent  $\alpha_J$  of a power law fit for  $dN_{sub}/dJ \sim J^{\alpha_J}$  in the regime  $J < 10^{17} \text{ GeV}^2 \text{ cm}^{-5}$ . The fitting uncertainty is  $\sigma_{\alpha_J} \lesssim 5 \cdot 10^{-2}$  for all models.

<sup>1</sup> Subhalos are drawn by CLUMPY down to a user-defined accuracy in terms of the  $J$ -factor contribution from subhalos compared to the background emission from the unclustered “smooth” DM in the halo. For all models, all subhalos yielding  $J \gtrsim 10^{16} \text{ GeV}^2 \text{ cm}^{-5}$  are correctly simulated in this study.

### 3. Setup of the simulations

The Galactic DM halo is simulated for a circular field of view (FOV) centered at the Galactic pole  $b = 90^\circ$  with diameter  $d = 120^\circ$ . This FOV covers a fraction  $f_{\text{sky}} = 25\%$  of the sky. The choice of the FOV is motivated by the projected extragalactic sky survey achievable with CTA [15], an area where it can be searched also for Galactic DM subhalos. To take into account the statistical variation of the observed sky in terms of cosmic variance, each simulation is rerun 10.000 times. In the same way, the search for DM subhalos within a single FOV of  $d = 10^\circ$  is investigated, using a separate simulation.

### 4. Results

In Fig. 2, the average integrated *source count distribution* is shown, defined as

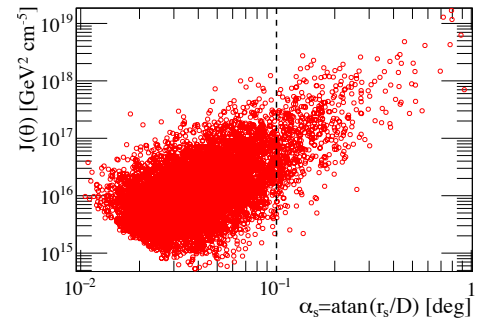
$$N_{\text{sub}}(\geq J) = \int_J^\infty \frac{dN_{\text{sub}}}{dJ'} dJ'. \quad (4.1)$$

Here,  $N_{\text{sub}}(\geq J)$  represents the number of subhalos on the chosen patch of the sky brighter than a given  $J$ -factor. The DM subhalo source count distribution  $dN_{\text{sub}}/dJ$  can be described by a power-law according to  $dN_{\text{sub}}/dJ \sim J^{-\alpha_J}$  in the regime  $J < 10^{17} \text{ GeV}^2 \text{ cm}^{-5}$ . In Tab. 1, the fit results of the exponent  $\alpha_J$  are listed for each of the 16 models. These results indicate that the  $dN_{\text{sub}}/dJ$  distribution is dominantly determined by the mass distribution  $dN_{\text{sub}}/dM_{\text{sub}} \sim M_{\text{sub}}^{-\alpha_m}$ .

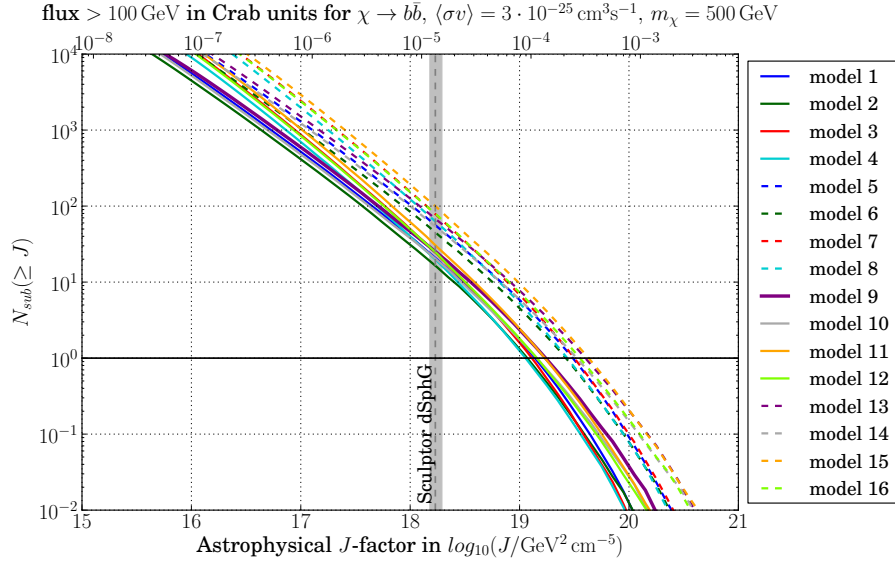
The integration angle, within which the  $J$ -factor of each subhalo is calculated, is set to  $\theta_1 = 0.1^\circ$  to illustrate the power-law behavior of the source count distribution. In Fig. 1, it is shown for model 9 that the largest part of most of the subhalos is encompassed within  $\theta_1$ . Similar results are obtained for all models. As the annihilation  $\gamma$ -ray emission is strongly peaked at the subhalo center, only the signal of the most extended subhalos gets significantly truncated by  $\theta_1$ . Fig. 1 suggests a correlation between the apparent size and the brightness of the subhalos. Therefore, a cut-off of the power-law distribution at high  $J$  is visible in Fig. 2. Continuous emission from the unclustered DM (“smooth halo”) and other background decreases proportional to  $\theta^2$ , faster than the subhalos’ flux [16]. Consequently, a smaller signal region might improve the signal-to-noise ratio in a real experiment.

Only a slight correlation between the subhalo masses and their  $J$ -factors has been found for all models (not shown). From this follows that distant massive dSphG-like subhalos and nearby low-mass subhalos similarly account for the highest  $J$ -factors.

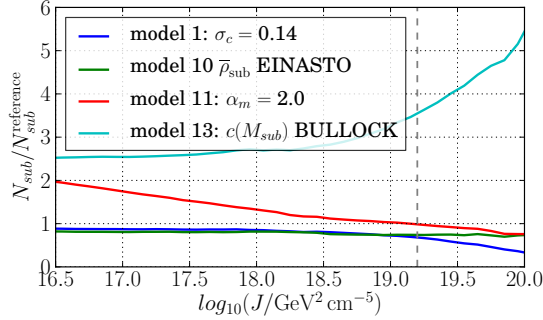
For a given DM annihilation channel, cross section and particle mass, the fluxes corresponding to the  $J$ -factors are plotted in Fig. 2. For the annihilation channel  $\chi\chi \rightarrow b\bar{b}$ , CTA will be most sensitive to a DM particle mass of  $500 \text{ GeV} \lesssim m_\chi \lesssim 1 \text{ TeV}$  [17]. For this mass range, the annihilation cross section is already constrained to  $\langle\sigma v\rangle \lesssim 3 \cdot 10^{-25} \text{ cm}^3 \text{ s}^{-1}$  by observations of the dSphGs by the *Fermi*-LAT instrument [18].



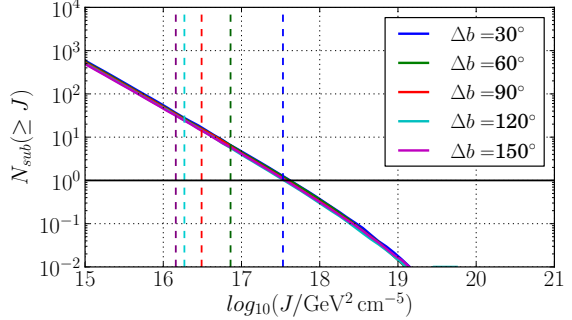
**Figure 1:** Relation between the angular size  $\alpha_s$  and the brightness  $J(\theta_1 = 0.1^\circ)$  of the subhalos for the survey region & model 9.  $D$  is the subhalo distance from the observer. The dashed line indicates  $\alpha_s = \theta_1$ .



**Figure 2:** Integrated source count distribution,  $N_{sub}$ , for all 16 benchmark models on  $f_{sky} = 0.25$  and for  $\theta_1 = 0.1^\circ$ . The dependence of  $N_{sub}$  on the  $J$ -factor is shown on the lower abscissa. On the upper abscissa, the dependence of  $N_{sub}$  on the flux is shown, for the assumed particle physics model. The continuous/dashed lines distinguish the SANCHEZ/BULLOCK concentration model. The corresponding  $J$ -factor from Sculptor dSphG for  $\theta_1 = 0.1^\circ$  [19] is plotted for comparison.



**Figure 3:**  $N_{sub}$  of the four models that differ only by the depicted parameter from model 9 ( $\sigma_c = 0.24$ ,  $\bar{\rho}_{sub}$  MADAU,  $\alpha_m = 1.9$ ,  $c(M_{sub})$  SANCHEZ), normalized to  $N_{sub}$  of model 9. The dashed line marks the  $J$  of the brightest subhalo expected ( $N_{sub}(\geq J) = 1$ ) for model 9.



**Figure 4:**  $N_{sub}$  for model 9 and a FOV diameter of  $d = 10^\circ$  at different distances  $\Delta b$  from the Galactic center. The maximum difference between all solid curves is  $< 10\%$ . The dashed lines denote the mean emission from the smooth galactic DM component for each FOV, and  $\theta_1 = 0.1^\circ$ .

Thus, in the range of the considered models and a coverage of  $f_{sky} = 25\%$ , a survey sensitivity to some  $10^{-4}$  of the Crab Nebula flux is necessary to discover at least one subhalo above 100 GeV with IACTs, provided that  $\langle\sigma v\rangle \approx 3 \cdot 10^{-25} \text{ cm}^3 \text{ s}^{-1}$ .

Fig. 3 shows the impact of the varied parameters on  $N_{sub}$  compared to the reference model. For the brightest subhalos, the model uncertainty is generally  $\lesssim 50\%$ . An exception from this occurs when switching from the SANCHEZ to the BULLOCK concentration model for  $c(M_{sub})$ . In this case, the number of brightest subhalos is increased by more than a factor of three.

Fig. 4 shows the simulation results for  $N_{sub}$  within a single CTA-like FOV of  $d = 10^\circ$ . Here,  $N_{sub}$  is given for different distances  $\Delta b$  from the Galactic center, perpendicular to the Galactic plane. Two conclusions can be drawn from this. First, Fig. 4 indicates that at  $r_\odot$ , the subhalos appear almost isotropically in the sky. Consequently, the relation  $N_{sub}(\Delta\Omega_1)/N_{sub}(\Delta\Omega_2) \approx \Delta\Omega_1/\Delta\Omega_2$  holds for arbitrary positions in the sky (and is confirmed by comparing Fig. 2 and 4). However, pointing close to the Galactic center heavily decreases the contrast between the subhalos and the smooth DM halo emission, the latter indicated by the dashed lines. Second, a sensitivity to  $J$ -factors of  $J \lesssim 10^{18} \text{ GeV}^2 \text{ cm}^{-5}$  is needed to resolve at least one subhalo within a given single FOV. Such  $J$ -factors are about a factor 10 smaller than those of the most promising dSphGs, where observation times of  $\gg 500 \text{ hrs}$  are expected to reach competitive limits  $\langle \sigma v \rangle \approx 10^{-25} \text{ cm}^3 \text{ s}^{-1}$  with CTA [17].

In Fig. 5, the *fluctuation angular power spectra* (fluctuation APS) of the subhalos are shown, comparing two of the four varied parameters against the reference model. The fluctuation APS is a measure of the relative fluctuation, or *small scale anisotropy*, of a signal, which can be possibly detected above a much larger background holding less fluctuation. The fluctuation APS  $C_\ell^F$  of a map  $I(\vartheta, \varphi)$ , defined on a patch  $f_{sky}$  of the sphere ( $I \equiv 0$  elsewhere) is given by

$$C_\ell^F = \frac{1}{f_{sky} \bar{I}^2 (2\ell + 1)} \sum_{m=-\ell}^{+\ell} |a_{\ell m}|^2, \quad (4.2)$$

with the spherical coordinates  $\vartheta$  and  $\varphi$ , and  $\bar{I}$  the mean value of  $I$  on the defined domain. The  $a_{\ell m}$  are the coefficients of the map decomposed into spherical harmonics  $Y_{\ell m}$ ,

$$I(\vartheta, \varphi) = \sum_{\ell=0}^{\ell_{max}} \sum_{m=-\ell}^{+\ell} a_{\ell m} Y_{\ell m}(\vartheta, \varphi). \quad (4.3)$$

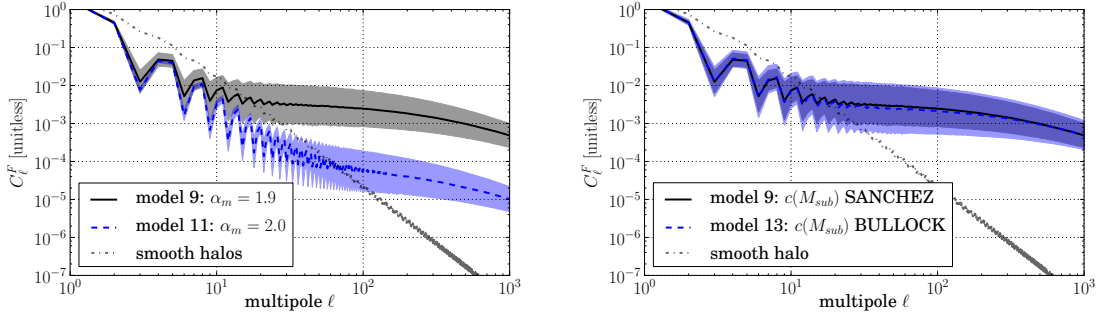
For Fig. 5, an integration angle  $\theta_2 = 0.015^\circ$  is chosen, which corresponds to half the value of the maximum resolution projected for CTA at energies above 1 TeV [15].

For a FOV with diameter  $d \lesssim 10^\circ$  and an angular resolution of  $\sigma_{PSF} \gtrsim 0.03^\circ$  [15], CTA could be sensitive to multipoles  $100 \lesssim \ell \lesssim 1000$  [6]. As shown in Fig. 5, in this regime, the fluctuation APS remains largely constant for all models, with only a mild lowering of the amplitudes  $C_\ell^F$  towards higher multipoles. Only in the case of the variation of  $\alpha_m$  from 1.9 to 2.0, a drastic change for the fluctuation APS results. This is shown in Fig. 5 (left). For  $\alpha_m = 2.0$ , a substantially higher amount of total mass is bound into subhalos. Although this results into a boost of the overall signal, the effective fluctuation heavily decreases because of the large background of the low-mass subhalos. For all other cases, the fluctuations of the subhalos remain largely the same. This applies even for the significant signal boost for the BULLOCK case, which is shown in Fig. 5 (right).

## 5. Conclusions

Different models of the abundance and properties of Galactic DM subhalos have been compared. Within the varied parameters, the source count distribution  $dN_{sub}/dJ$  of these subhalos has been constrained to an uncertainty band of one order of magnitude. The modeling of the inner concentration  $c(M_{sub})$  causes the biggest systematic uncertainty in the number of detectable subhalos. The subhalos appear nearly isotropically distributed in the sky, allowing a telescope to point as





**Figure 5:** Fluctuation APS  $C_\ell^F$  of the subhalos on  $f_{sky} = 0.25$ , comparing the variation of the index  $\alpha_m$  (left) and the  $c(M_{sub})$  model (right) against model 9. The fluctuation APS of the smooth halo on the simulated patch of the sky is plotted for comparison. The bands indicate the  $1\sigma$  containment ranges around the mean spectrum of the simulated galactic halo, assuming a log-normal variation of the  $C_\ell^F$ . Note that for the left panel, also the smooth contribution varies between the compared models. However, this difference in  $C_\ell^F$  between the smooth contributions is smaller than the resolution of the plots and not distinguished here.

far as possible away from the Galactic plane to reduce the background from the smooth DM halo emission and diffuse astrophysical foregrounds.

However, for all benchmark models, Galactic subhalos are probably too faint to be a promising target for current and next-generation IACTs, provided the projected available observation times. Still for a sensitivity to a fraction of  $10^{-4}$  of the Crab Nebula flux above 100 GeV, a detection of individual subhalos might be possible only for cross sections  $\langle\sigma v\rangle \gtrsim 10^{-25} \text{ cm}^3 \text{ s}^{-1}$ , which are already in tension with recent measurements. Pointed observations of the Galactic center and the dSphGs with CTA promise competitive limits on the annihilation cross section, probably requiring much less observation time.

At the same time, it has been shown that a source count distribution with  $\alpha_J \lesssim 2.1$  results in a fluctuation APS up to the order of  $10^{-3} \lesssim C_\ell^F \lesssim 10^{-2}$  in the multipole range of interest of Cherenkov telescopes. In principle, such a comparably high fluctuation might distinguish dark matter from other, more dominant unresolved source classes. Possibly, even more competitive limits on the annihilation cross section can be obtained than by the search for individual subhalos [6]. However, given the low intensity of the overall signal from Galactic DM subhalos, these limits are most likely still not competitive with those obtainable from (combined) pointed observations.

We did not assess the extragalactic contribution to DM induced fluctuations of the IGRB. However, the results of [8] indicate that the extragalactic emission is expected to be similar to the Galactic signal in magnitude and in angular scales, suggesting a similar sensitivity to IGRB fluctuations as for the Galactic subhalo case. Still, it should be emphasized that all indirect DM targets suffer from major systematic uncertainties. With this in mind, the search for fluctuations in the IGRB should not be neglected as a complementary strategy to unveil the nature of DM.

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